

Analysis of Impulsive Response of Mechanical and Electrical Oscillators by Rohit Transform

Rohit Gupta^{*1}, Rahul Gupta^{*2}, Ajay Kumar³

¹Lecturer of Physics, Department of Applied Sciences, Yogananda College of Engineering and Technology, Jammu (J&K), India

²Lecturer of Physics, Department of Physics, SDMP Public Hr. Secondary School, Karwanda Balwal,, Jammu (J&K), India

³Assistant Professor, Department of Physics, Govt. Degree College, Kishtwar (J&K), India

*Corresponding author E-mail id: guptara702@gmail.com

Abstract — The analysis of impulsive responses of mechanical and electrical oscillators is a well-known problem in applied sciences and engineering. In this paper, the impulsive responses of mechanical and electrical oscillators are obtained by the new integral transform ‘Rohit transform’. This paper brings up the Rohit transform as a new technological approach for obtaining the impulsive responses of mechanical and electrical oscillators. The results obtained by applying the Rohit transform are the same as obtained with other integral transforms or approaches.

Keywords — Mechanical and Electrical Oscillators; Impulsive Responses; Rohit Transform.

1. Introduction

The mechanical and electrical oscillators are generally analyzed by different integral transforms and approaches [1]-[12]. This paper analyses the mechanical and electrical oscillators by new integral transform ‘Rohit transform’ to obtain their impulsive responses. The Rohit Transform is a new integral transform which has been proposed by the author Rohit Gupta in the recently in the year 2020 and generally, it has been applied to boundary value problems in most of the science and engineering disciplines such as for the analysis of radioactive decay problem [13], for solving the Schrodinger equation [14], for the Analysis of RLC circuits with exponential excitation sources [15], for the Analysis of Basic Series Inverter [16], for the analysis of Electric Network Circuits with Sinusoidal Potential Sources [17], for the analysis of uniform infinite fin [18], Response of RLC network circuit with steady source [19], for the analysis of one-way streamline flow between parallel plates [20], Response of a Permanent Magnet Moving Coil Instrument [21], Heat Conducted Through Fins Of Varying Cross-Sections [22], Response of an Undamped Forced Oscillator [23], Analysis Of Damped Mechanical And Electrical Oscillators [24]. The purpose of paper is to demonstrate the Rohit transform for analyzing the impulsive responses of mechanical and electrical oscillators. The energy of a simple harmonic oscillator persists indefinitely without the reduction of amplitude and remains constant throughout the motion. In practice, the oscillator is acted upon by the frictional forces arising from the viscosity of the medium or from within the system itself which causes the reduction of amplitude and hence the energy of the oscillator with time [1], [10]-[12]. In electrical oscillator, energy is dissipated due to the presence of resistance of the circuit [7]-[9]. However, the presence of small damping force does not have any significant effect on

the undamped oscillations of the oscillator. In such a case, the frictional forces acting on the oscillator are directly proportional to the velocity of the oscillator only [1]-[12].

The Rohit transform of $g(y)$, $y \geq 0$ is denoted by $G(r)$ and is given by,

$$G(r) = r^3 \int_0^{\infty} e^{-ry} g(y) dy,$$

provided the integral is convergent, where r may be a real or complex parameter [14]-[17].

The RT of some elementary functions are

- ❖ $R\{y^n\} = \frac{n!}{r^{n-2}}$, where $n = 0, 1, 2, 3 \dots$
- ❖ $R\{e^{by}\} = \frac{r^3}{r-b}$, $r > b$
- ❖ $R\{\sin by\} = \frac{br^3}{r^2+b^2}$, $r > 0$
- ❖ $R\{\sinh by\} = \frac{br^3}{r^2-b^2}$, $r > |b|$
- ❖ $R\{\cos by\} = \frac{r^4}{r^2+b^2}$, $r > 0$
- ❖ $R\{\cosh by\} = \frac{r^4}{r^2-b^2}$, $r > |b|$
- ❖ $R\{\delta(y)\} = r^3$, where $\delta(y)$ is Dirac delta function.

1.1 Some Properties Rohit Transform

$$\{\delta(y)\} = r^3 \int_0^{\infty} e^{-ry} \delta(y) dy$$

Since the Dirac delta function [1], [12] $\delta(y)$, is defined as $\delta(y) = 0$ when $y \neq 0$ and $\int_{-\infty}^{\infty} g(y) \delta(y) dy = g(0)$, therefore, the above integral can be rewritten as

$$R\{\delta(y)\} = r^3 e^{-r(0)}$$

$$R\{\delta(y)\} = r^3$$

$$\text{Hence } R\{\delta(y)\} = r^3$$

Some properties of the Rohit transform on partial derivatives are discussed as below:

The Rohit Transform (RT) of first derivative of $g(y)$ i.e. RT of $\frac{\partial g(y)}{\partial y}$ w.r.t. y , ($y \geq 0$) [14]-[16] is given by

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = r^3 \int_0^\infty e^{-ry} \frac{\partial g(y)}{\partial y} dy$$

Integrating by parts and applying limits, we get

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = r^3 \left\{ -g(0) - \int_0^\infty -r e^{-ry} g(y) dy \right\}$$

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = r^3 \left\{ -g(0) + r \int_0^\infty e^{-ry} g(y) dy \right\}$$

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = r^3 \{-g(0) + rR\{g(y)\}\}$$

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = rG(r) - r^3 g(0)$$

Hence $R \left\{ \frac{\partial g(y)}{\partial y} \right\} = rG(r) - r^3 g(0)$

Since $R \left\{ \frac{\partial g(y)}{\partial y} \right\} = rR\{g(y)\} - r^3 g(0)$,

therefore, on replacing $g(y)$ by $\frac{\partial g(y)}{\partial y}$ and $\frac{\partial g(y)}{\partial y}$ by $\frac{\partial^2 g(y)}{\partial y^2}$,

we have

$$R \left\{ \frac{\partial^2 g(y)}{\partial y^2} \right\} = rR \left\{ \frac{\partial g(y)}{\partial y} \right\} - r^3 \frac{\partial g(0)}{\partial y}$$

$$R \left\{ \frac{\partial^2 g(y)}{\partial y^2} \right\} = r\{rR\{g(y)\} - r^3 g(0)\} - r^3 \frac{\partial g(0)}{\partial y}$$

$$R \left\{ \frac{\partial^2 g(y)}{\partial y^2} \right\} = r^2 R\{g(y)\} - r^4 g(0) - r^3 \frac{\partial g(0)}{\partial y}$$

$$R \left\{ \frac{\partial^2 g(y)}{\partial y^2} \right\} = r^2 G(r) - r^4 g(0) - r^3 \frac{\partial g(0)}{\partial y}$$

Hence $R \left\{ \frac{\partial^2 g(y)}{\partial y^2} \right\} = r^2 G(r) - r^4 g(0) - r^3 \frac{\partial g(0)}{\partial y}$

Thus the RT of derivatives of $g(y)$ with respect to y ($y \geq 0$) are given by

$$R \left\{ \frac{\partial g(y)}{\partial y} \right\} = rG(r) - r^3 g(0),$$

$$R \left\{ \frac{\partial^2 f(y)}{\partial y^2} \right\} = r^2 G(r) - r^4 g(0) - r^3 g'(0),$$

$$R \left\{ \frac{\partial^3 f(y)}{\partial y^3} \right\} = r^3 G(r) - r^5 g(0) - r^4 g'(0) - r^3 g''(0),$$

And so on.

2. Material and Method

2.1 Mechanical (Damped) Oscillator

The differential equation of damped oscillator subjected to the impulsive force [1], [3], [12] is given by $\ddot{x}(t) + 2\alpha\dot{x}(t) + \omega^2 x(t) = g(t)$ (1), where $g(t)$

= $\gamma\delta(t)$, $\delta(t)$ is a delta function (impulsive force per unit mass) and γ is its strength, $2\alpha = \frac{r}{m}$ represents the damping

constant per unit mass, $\omega = \sqrt{\frac{k}{m}}$ represents the natural frequency of the oscillator. For a lightly damped oscillator, $\alpha < \omega$.

The initial boundary conditions as follows [3]-[5], [25], [26]:

- i. If we measure the time from the instant when the oscillator is crossing its mean position, then at $t = 0$, $x(0) = 0$.
- ii. Also, at the instant $t = 0^+$ (i.e. just after applying the impulsive force), we assume that the velocity of the oscillator is maximum i.e. $\dot{x}(0^+) = v_0$.

The Rohit transform of (1) provides

$$q^2 \bar{x}(q) - q^4 x(0) - q^3 \dot{x}(0) + 2\alpha\{q\bar{x}(q) - q^3 x(0)\} + \omega^2 \bar{x}(q) = \gamma q^3 \dots (2)$$

Here $\bar{x}(q)$ denotes the Rohit transform of $x(t)$.

Applying boundary conditions $x(0) = 0$ and $\dot{x}(0) = v_0$ and simplifying (2), we get

$$\bar{x}(q) = \frac{(\gamma + v_0)q^3}{q^2 + 2\alpha q + \omega^2}$$

Or

$$\bar{x}(q) = \frac{(\gamma + v_0)q^3}{(q + \beta_1)(q + \beta_2)}$$

where $\beta_1 = \alpha + i\sqrt{\omega^2 - \alpha^2}$ and $\beta_2 = \alpha - i\sqrt{\omega^2 - \alpha^2}$ such that $\beta_1 - \beta_2 = 2i\sqrt{\omega^2 - \alpha^2}$

$$\text{Or } \bar{x}(q) = \frac{(\gamma + v_0)q^3}{(\beta_2 - \beta_1)(q + \beta_1)} - \frac{(\gamma + v_0)q^3}{(\beta_2 - \beta_1)(q + \beta_2)}$$

Applying inverse Rohit Transform, we get

$$x(t) = (\gamma + v_0) \frac{[e^{-\beta_1 t} - e^{-\beta_2 t}]}{(\beta_2 - \beta_1)} \dots \dots \dots (3)$$

To find γ , applying condition $\dot{x}(0) = v_0$ and solving, we get $x = 0 \dots \dots \dots (4)$

Using (4) in (3), we get

$$x(t) = v_0 \frac{[e^{-\beta_1 t} - e^{-\beta_2 t}]}{(\beta_2 - \beta_1)}$$

Or

$$x(t) = v_0 e^{-\alpha t} \frac{[e^{i\sqrt{\omega^2 - \alpha^2} t} - e^{-i\sqrt{\omega^2 - \alpha^2} t}]}{2i\sqrt{\omega^2 - \alpha^2}}$$

Or

$$x(t) = \frac{v_0 e^{-\alpha t}}{\sqrt{\omega^2 - \alpha^2}} \sin \sqrt{\omega^2 - \alpha^2} t \dots (5)$$

When $r = 0$, $\alpha = \frac{r}{2m} = 0$, then (5) reduces to

$$x(t) = \frac{v_0}{\omega} \sin \omega t$$

The equation (5) provides the impulsive response of a lightly damped oscillator and reveals that it is independent

on the strength of the impulsive force. Also, the behavior of oscillator is oscillatory with amplitude of oscillations reduces with time exponentially. When the damping force is zero, the amplitude of oscillations remains constant.

For an overdamped oscillator [5], $\alpha > \omega$, therefore, replacing $\sqrt{\omega^2 - \alpha^2}$ by $i\sqrt{\alpha^2 - \omega^2}$ in (5), the displacement of an overdamped oscillator is given by

$$x(t) = \frac{v_0 e^{-\alpha t}}{i\sqrt{\alpha^2 - \omega^2}} \sin i\sqrt{\alpha^2 - \omega^2} t$$

Or

$$x(t) = \frac{v_0 e^{-\alpha t}}{\sqrt{\alpha^2 - \omega^2}} \sinh \sqrt{\alpha^2 - \omega^2} t \dots (6)$$

This equation (6) provides the impulsive response of heavily damped oscillator and reveals that the motion of heavily damped oscillator is non-oscillatory.

2.2 Electrical (Damped) Oscillator

The differential equation of electrical oscillator (LRC circuit) subjected to the impulsive Potential [8], [9], [12] is given by

$$Q(t) + 2\alpha \dot{Q}(t) + \omega^2 Q(t) = g(t) \dots (7), \text{ where } g(t)$$

$= \gamma \delta(t)$, where $g(t) = \gamma \delta(t)$, $\delta(t)$ is a delta potential

(impulse) and γ is its strength, $\omega = \sqrt{\frac{1}{LC}}$ represents the

angular frequency of the electrical oscillator, $2\alpha = \frac{R}{L}$

represents the damping coefficient. $Q(t)$ is the

instantaneous charge.

The initial boundary conditions [10], [11] as follows:

- (i) At $t = 0$, $Q(0) = 0$.
- (ii) Also, at the instant $t = 0^+$ (i.e. just after applying the impulsive potential, we assume that the current in the circuit is maximum i.e. $\dot{Q}(0^+) = i_0$).

The Rohit transform of (1) provides

$$q^2 \bar{Q}(q) - q^4 Q(0) - q^3 \dot{Q}(0) + 2\alpha \{q \bar{Q}(q) - q^3 Q(0)\} + \omega^2 \bar{Q}(q) = \gamma q^3 \dots (8)$$

Here $\bar{Q}(q)$ denotes the Rohit transform of $Q(t)$.

Applying boundary conditions [12] $Q(0) = 0$

and $\dot{Q}(0) = i_0$ and simplifying (8), we get

$$\bar{Q}(q) = \frac{(\gamma + i_0)q^3}{q^2 + 2\alpha q + \omega^2}$$

Or

$$\bar{Q}(q) = \frac{(\gamma + i_0)q^3}{(q + \beta_1)(q + \beta_2)}$$

where $\beta_1 = \alpha + i\sqrt{\omega^2 - \alpha^2}$ and $\beta_2 = \alpha - i\sqrt{\omega^2 - \alpha^2}$ such that $\beta_1 - \beta_2 = 2i\sqrt{\omega^2 - \alpha^2}$.

$$\text{Or } \bar{Q}(q) = \frac{(\gamma + i_0)q^3}{(\beta_2 - \beta_1)(q + \beta_1)} - \frac{(\gamma + i_0)q^3}{(\beta_2 - \beta_1)(q + \beta_2)}$$

Applying inverse Rohit Transform, we get

$$Q(t) = (\gamma + i_0) \frac{[e^{-\beta_1 t} - e^{-\beta_2 t}]}{(\beta_2 - \beta_1)} \dots (9)$$

To find γ , applying condition $Q(0) = i_0$ and solving, we get $\gamma = 0 \dots \dots (10)$

Using (10) in (9), we get

$$Q(t) = i_0 \frac{[e^{-\beta_1 t} - e^{-\beta_2 t}]}{(\beta_2 - \beta_1)}$$

Or

$$Q(t) = i_0 e^{-\alpha t} \frac{[e^{i\sqrt{\omega^2 - \alpha^2} t} - e^{-i\sqrt{\omega^2 - \alpha^2} t}]}{2i\sqrt{\omega^2 - \alpha^2}}$$

Or

$$Q(t) = \frac{i_0 e^{-\alpha t}}{\sqrt{\omega^2 - \alpha^2}} \sin \sqrt{\omega^2 - \alpha^2} t \dots (11)$$

When $R = 0$, $\alpha = \frac{R}{2L} = 0$, then (11) reduces to

$$Q(t) = \frac{i_0}{\omega} \sin \omega t$$

The equation (11) provides the impulsive response of an electrical oscillator and reveals that it is independent on the strength of the impulsive potential. Also, the behavior of oscillator (charge) is oscillatory with the amplitude of oscillations reduces with time exponentially. The decrease in amplitude i.e. damping depends upon resistance R in the circuit. Such a damping is called resistance damping [2] [9] [12]. If $R = 0$, the amplitude would remain constant. Hence in LRC circuit, the resistance is the only dissipative element.

3. Conclusion

In this paper, we have successfully applied the Rohit Transform to obtain the impulsive responses of damped mechanical and electrical oscillators and exemplified it for successfully analyzing the theory of damped mechanical and electrical oscillators. A new and different method is exploited for obtaining the impulsive responses of damped mechanical and electrical oscillators.

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